

Analytic moduli for unfoldings of germs of generic analytic diffeomorphisms with a codimension k parabolic point*

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Abstract

In this paper we provide a complete modulus of analytic classification for germs of generic analytic families of diffeomorphisms which unfold a parabolic fixed point of codimension k . We start by showing that a generic family can be “prepared”, i.e. brought to a prenormal form $f_\epsilon(z)$ in which the multi-parameter ϵ is almost canonical (up to an action of $\mathbb{Z}/k\mathbb{Z}$). As in the codimension 1 case treated in [7], we show that the Ecalle-Voronin modulus can be unfolded to give a complete modulus for such germs. For this purpose, we define unfolded sectors in z -space that constitute natural domains on which the map f_ϵ can be brought to normal form in an almost unique way. The comparison of these normalizing changes of coordinates on the different sectors forms the analytic part of the modulus. This construction is performed on sectors in the multi-parameter space ϵ such that the closure of their union provides a neighborhood of the origin in parameter space.

1 Introduction

This paper is part of a large program to classify and understand the finite codimension singularities of analytic dynamical systems and their unfoldings. The first tool for this purpose is to look at normal forms. A certain number of singularities have normal forms depending on a finite number of parameters which are formal invariants of the system. In a series of papers ([7], [1], [11], [13], [12], [17], [6]) we have developed the program for singularities with common features: their normal form have a finite number of formal parameters, called *formal invariants*, and there exist formal changes of coordinates to normal form which are divergent and k -summable in some sense. The singularities studied come from the merging of $k + 1$ simple singularities ([7], [1], [12], [6] when $k = 1$, and [17] for a codimension k saddle-node of a planar vector field), or from the merging of a singularity with a periodic orbit ([11] and [13]). Examples of such singularities are the fixed points of a 1-dimensional diffeomorphism with resonant multiplier (the multiplier is a root of unity), the saddle-nodes, the resonant saddles (the quotient of eigenvalues is a negative rational number) or the weak foci of a 2-dimensional vector field, and the irregular non resonant singularities of a linear differential system.

The meaning of the formal invariants is best understood when one unfolds the system. Indeed, in all cases except [6], there are natural *canonical parameters* which are analytic

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invariants for the family. The formal invariants unfold analytically in the parameters. In all generic cases (simple singularities), the sum of the number of parameters and of the number of formal invariants is exactly equal to the sum of the number of (linear) invariants at each simple singular point in the unfolding (multipliers or eigenvalues).

The analytic classification of germs of generic analytic families unfolding these singularities is done by giving a complete modulus of analytic classification for each family. The modulus comprises a formal part described above, and an analytic part which is given by the unfolding of the analytic part of the modulus of the singularity itself. To describe the modulus, we cover a fixed neighborhood of the singularity with $2k$ sectors depending on the parameters. On each sector, we construct normalizing changes of coordinates to the normal form. The sectors are such that the normalizing changes of coordinates are almost unique (up to symmetries of the normal form on the full neighborhood). The comparison of the normalizing changes of coordinates on the intersections of the sectors provides the modulus. The construction can be performed depending analytically on the parameters on some sectors in parameter space, the closure of which cover a full neighborhood of the origin in parameter space. The analytic part of the modulus of the singularity is the limit of the analytic part of the modulus of the unfolding: this allows to interpret the modulus of the codimension k singularity and to understand why it is so complex. Indeed, there are rigid models at each unfolded singularity which are glued in a non trivial way. The modulus at the limit is the limit of these non trivial gluings. Also the unfolded singularities can be resonant for sequences of parameter values converging to the origin, in which case they have themselves formal invariants. The limit of these formal invariants are also encoded in the analytic part of the modulus of the limit singularity: this is the *parametric resurgence phenomenon* discussed in Section 6.

In this paper we build a complete modulus of analytic classification for germs of generic analytic diffeomorphisms with a codimension k parabolic point. As a first step, in Section 2, we show that the family can be *prepared*, i.e. brought through a change of variable and parameter to a form in which the parameters are analytic invariants up to an action of $\mathbb{Z}/k\mathbb{Z}$. In Section 4, we then describe the sectors in the variable and in the parameter space, and we construct the Fatou coordinates which normalize the diffeomorphism on the different sectors. The comparison of the Fatou coordinates yields the analytic part of the modulus. This analytic part depends analytically on the parameters on each sector in parameter space. In Section 5, we then show that the modulus we have defined is a complete modulus of analytic classification, namely that two families with the same modulus are analytically conjugate, with the conjugacy depending analytically on the parameters. In Section 6, we briefly discuss the interpretation of the modulus in the parametric resurgence phenomenon.

The next step in this program is to identify precisely the set of realizable moduli. For this purpose, it is needed to identify the sufficient conditions for a modulus to be realized as the modulus of a germ of analytic family unfolding a diffeomorphism with a parabolic point of codimension k . We hope to address this question in a forthcoming paper.

2 Preparation of the family

Definition 2.1 *A germ of an analytic k -parameter family f_η unfolding a diffeomorphism of the form*

$$f_0(z) = z + z^{k+1} + o(z^{k+1}), \quad (2.1)$$

can be written (modulo a translation) in the form

$$f_\eta(z) = z + p_\eta(z)(1 + O(|(z, \eta)|)),$$

where

$$p_\eta(z) = z^{k+1} + c_{k-1}(\eta)z^{k-1} + \dots c_1(\eta)z + c_0(\eta).$$

It is called a germ of generic analytic k -parameter family if the mapping $\eta \mapsto (c_0(\eta), \dots, c_{k-1}(\eta))$ is a germ of analytic diffeomorphism from $(\mathbb{C}^k, 0)$ to $(\mathbb{C}^k, 0)$.

Theorem 2.2 *For any germ of generic analytic k -parameter family f_η unfolding a diffeomorphism of the form (2.1) there exists a germ of change of coordinate and parameter to a prepared form*

$$\bar{f}_\epsilon(Z) = Z + P_\epsilon(Z)(1 + Q_\epsilon(Z) + P_\epsilon(Z)R_\epsilon(Z)) \quad (2.2)$$

with

$$P_\epsilon(z) = Z^{k+1} + \epsilon_{k-1}Z^{k-1} + \dots + \epsilon_1Z + \epsilon_0, \quad (2.3)$$

such that:

- $Q_\epsilon(Z) = c_k(\epsilon)Z^k + \dots + c_1(\epsilon)Z + c_0(\epsilon)$, with $c_0(0) = \dots = c_{k-1}(0) = 0$;
- If Z_1, \dots, Z_{k+1} are the fixed points of \bar{f}_ϵ , if $\lambda_j(\epsilon)$ is defined by $\lambda_j(\epsilon) = \bar{f}'_\epsilon(Z_j)$, if the germ of analytic function $a(\epsilon)$ is defined through

$$\sum 1/\ln(\lambda_j(\epsilon)) = a(\epsilon) \quad (2.4)$$

when the Z_j are distinct, and the vector field v_ϵ is defined by

$$v_\epsilon(Z) = \frac{P_\epsilon(Z)}{1 + a(\epsilon)Z^k} \frac{\partial}{\partial Z}, \quad (2.5)$$

then

$$\ln \lambda_j(\epsilon) = \ln(\bar{f}'_\epsilon(Z_j)) = v'_\epsilon(Z_j). \quad (2.6)$$

Moreover, the parameters ϵ are almost unique in the following sense. If $\tilde{f}_\epsilon(\tilde{Z})$ is another prepared form for the same unfolding f_η , then there exists τ such that $\tau^k = 1$ and

$$(\tilde{Z}, \tilde{\epsilon}_{k-1}, \dots, \tilde{\epsilon}_0) = (\tau Z + o(Z), \tau^2 \epsilon_{k-1}, \tau^3 \epsilon_{k-2}, \dots, \epsilon_1, \tau \epsilon_0).$$

PROOF. Using a translation in z and a change of parameters, we can suppose that the fixed points of f_η are given by the zeros of

$$p_\eta(z) = z^{k+1} + \eta_{k-1}z^{k-1} + \dots + \eta_1z + \eta_0.$$

Hence, we can write the unfolding in the form

$$f_\eta(z) = z + p_\eta(z)(1 + m_\eta(z)),$$

with $m_\eta(z) = O(\eta, z)$. Using the Weierstrass division theorem on m_η allows to write f_η in the form

$$f_\eta(z) = z + p_\eta(z)(1 + q_\eta(z) + p_\eta(z)n_\eta(z)) \quad (2.7)$$

with $q_\eta(z)$ a polynomial of degree k .

Let z_1, \dots, z_{k+1} be the fixed points of f_η . There exists a polynomial $S_\eta(z)$ of degree $\leq k$ such that at the points z_j we have

$$\ln f'_\eta(z_j) = p'_\eta(z_j)(1 + S_\eta(z_j)). \quad (2.8)$$

Such a polynomial is found by the Lagrange interpolation formula using

$$S_\eta(z_j) = \frac{\ln(1 + p'_\eta(z_j)(1 + q_\eta(z_j)))}{p'_\eta(z_j)} - 1.$$

Indeed, the formula works for distinct z_j . The limit exists when two fixed points coalesce (codimension 1 case). We can fill in for the other values of η by Hartogs's Theorem.

By Kostov's theorem [5], there exists a local analytic change of coordinate and parameters transforming the vector field

$$v_{1,\eta}(z) = p_\eta(z)(1 + S_\eta(z)) \frac{\partial}{\partial z}$$

into

$$v_\epsilon(Z) = P_\epsilon(Z)/(1 + a(\epsilon)Z^k) \frac{\partial}{\partial Z},$$

where P_ϵ is defined in (2.3). We apply this change of coordinate and parameter to f_η .

Claim: this brings f_η to a prepared form \bar{f}_ϵ . Indeed: it sends the zeros z_j of $p_\eta(z)$ to the zeroes Z_j of $P_\epsilon(Z)$. Since the z_j are the fixed points of f_η , their images are the fixed points Z_j of \bar{f}_ϵ . So

$$\bar{f}_\epsilon(Z) = Z + P_\epsilon(Z)G_\epsilon(Z) = Z + P_\epsilon(Z)(1 + Q_\epsilon(Z) + P_\epsilon(Z)L_\epsilon(Z)).$$

Let Z_j be a fixed point of \bar{f}_ϵ . Then, from (2.8),

$$f'_\eta(z_j) = \exp(p'_\eta(z_j)(1 + S_\eta(z_j))) = \exp(v'_{1,\eta}(z_j)).$$

Since multipliers at fixed points of diffeomorphisms (resp. eigenvalues at singular points of vector fields) are invariant under change of coordinates, then we have

$$\bar{f}'_\epsilon(Z_j) = \lambda_j(\epsilon) = f'_\eta(z_j) = \exp(v'_{1,\eta}(z_j)) = \exp(v'_\epsilon(Z_j))$$

which is what we need for a prepared family. Moreover, we see that

$$\bar{f}'_\epsilon(Z_j) = 1 + P'_\epsilon(Z_j)(1 + Q_\epsilon(Z_j))$$

can be calculated from P_ϵ and Q_ϵ alone.

We now need to show that $Q_0(Z) = O(Z^k)$. So we consider the particular case where $\epsilon_j = 0$ for $j \neq 1$. The fixed points are then $Z_0 = 0$, and Z_1, \dots, Z_k , with $Z_j^k + \epsilon_1 = 0$. This yields $P'_\epsilon(0) = \epsilon_1$ and $P_\epsilon(Z_j) = -k\epsilon_1$ for $j > 1$. We need have on one side $\bar{f}'_\epsilon(Z_j) = \exp\left(\frac{-k\epsilon_1}{1-a(\epsilon)\epsilon_1}\right)$, which yields $\bar{f}'_\epsilon(Z_j) = 1 - k\epsilon_1 + O(\epsilon_1^2)$. On the other side, we also have $\bar{f}'_\epsilon(Z_j) = 1 - k\epsilon_1(1 + Q_\epsilon(Z_j))$. Hence, if the coefficient $c_\ell(\epsilon)$ of Z^ℓ in Q_ϵ was such that $c_\ell(0) \neq 0$ for

$\ell < k$ and ℓ was minimum with this property, then this would yield $\bar{f}'_\epsilon(Z_j) = 1 - k\epsilon_1 + C\epsilon_1^{1+\frac{\ell}{k}}$ with $C \neq 0$, hence a contradiction.

We now show that the parameters are almost canonical. For that purpose, let us consider a prepared family, which we denote $f_\epsilon(z)$ (forgetting the bar). We will prove that the only freedom we have in the choice of a prepared family is a change

$$(z, \epsilon_{k-1}, \epsilon_{k-2}, \dots, \epsilon_0) \mapsto (\tau z + o(z), \tau^{-2}\epsilon_{k-1}, \tau^{-3}\epsilon_{k-2}, \dots, \tau\epsilon_0), \quad (2.9)$$

where $\tau = \exp(2\pi i \ell / k)$, $\ell = 0, \dots, k-1$. Indeed, let us suppose that two prepared families $f_\epsilon(z)$ and $\tilde{f}_{\tilde{\epsilon}}(\tilde{z})$ are conjugate under a map $(\tilde{\epsilon}, \tilde{z}) = (g(\epsilon), h_\epsilon(z))$: $\tilde{f}_{g(\epsilon)} = h_\epsilon \circ f_\epsilon \circ h_\epsilon^{-1}$. The fixed points z_j (resp. \tilde{z}_j) of the family f_ϵ (resp. $\tilde{f}_{\tilde{\epsilon}}$) are the same as the singular points of a vector field $P_\epsilon(z)/(1 + a(\epsilon)z^k) \frac{\partial}{\partial z} = v_\epsilon(z)$ (resp. $\tilde{P}_{\tilde{\epsilon}}(\tilde{z})/(1 + \tilde{a}(\tilde{\epsilon})\tilde{z}^k) \frac{\partial}{\partial \tilde{z}} = \tilde{v}_{\tilde{\epsilon}}(\tilde{z})$). Note that $a(\epsilon) = \tilde{a}(g(\epsilon))$. Then h_ϵ sends the fixed points z_j to the fixed points \tilde{z}_j . Hence, $h_\epsilon^*(\tilde{v}_{g(\epsilon)})(z) = P_\epsilon(z)U_\epsilon(z) \frac{\partial}{\partial z} = \bar{v}_\epsilon(z)$, where $U_\epsilon \neq 0$. Moreover, the vector fields v_ϵ and \bar{v}_ϵ have the same singular points with same eigenvalues. Hence, we can write

$$\bar{v}_\epsilon = P_\epsilon(z) \left(\frac{1}{1 + a(\epsilon)z^k} + P_\epsilon(z)M_\epsilon(z) \right) \frac{\partial}{\partial z} = v_\epsilon(1 + P_\epsilon(z)N_\epsilon(z)) \frac{\partial}{\partial z}.$$

There exists a change of coordinate K_ϵ such that $K_\epsilon^*(v_\epsilon) = \bar{v}_\epsilon$. Such a change of coordinate is given by the flow of v_ϵ under some time T_ϵ (see for instance Corollary 8.2 of [17]): $K_\epsilon = \Phi_{v_\epsilon}^{T_\epsilon}$, where T_ϵ is solution of

$$v_\epsilon(T_\epsilon) = -\frac{P_\epsilon(z)N_\epsilon(z)}{1 + P_\epsilon(z)N_\epsilon(z)}$$

which obviously has an analytic solution. Then $K_\epsilon^{-1} \circ h_\epsilon^*(\tilde{v}_{g(\epsilon)}) = v_\epsilon$. The result follows from the following theorem proved in [17]. \square

Theorem 2.3 [17] *We consider two vector fields, $v_\epsilon = P_\epsilon(z)/(1 + a(\epsilon)z^k) \frac{\partial}{\partial z}$ and $\tilde{v}_{\tilde{\epsilon}} = \tilde{P}_{\tilde{\epsilon}}(\tilde{z})/(1 + \tilde{a}(\tilde{\epsilon})\tilde{z}^k) \frac{\partial}{\partial \tilde{z}}$, which are conjugate in a neighborhood of the origin containing all the singular points. Then, necessarily the conjugacy between v_ϵ and $\tilde{v}_{\tilde{\epsilon}}$ is of the form (2.9) with τ some k -th root of unity.*

Remark 2.4 1. The meaning of the formal invariant a now becomes clear. The polynomial P_ϵ has $k+1$ roots and depends on k parameters only. An additional parameter is needed so that the eigenvalues at the singular points of v_ϵ are independent.

2. In a prepared family the parameters are almost canonical. This simplifies greatly the problem of classifying germs of families under analytic equivalence. Indeed, in general, two prepared families $f_\epsilon(z)$ and $\tilde{f}_{\tilde{\epsilon}}(\tilde{z})$ are analytically conjugate under an analytic diffeomorphism $(\tilde{\epsilon}, \tilde{z}) = (g(\epsilon), h_\epsilon(z))$ such that $\tilde{f}_{g(\epsilon)} = h_\epsilon \circ f_\epsilon \circ h_\epsilon^{-1}$. Here, we can limit ourselves to maps $g(\epsilon) = (\tau^{-2}\epsilon_{k-1}, \tau^{-3}\epsilon_{k-2}, \dots, \tau\epsilon_0)$ for τ some k -th root of unity. Note also that the parameters of the family of diffeomorphisms f_ϵ are the same as those of the vector field v_ϵ .

3 Some preliminaries and notations

The vector field $w_\epsilon = P_\epsilon(z) \frac{\partial}{\partial z}$ is invariant under

$$(z, t, \epsilon_{k-1}, \dots, \epsilon_0) \mapsto (cz, c^{-k}t, c^2\epsilon_{k-1}, \dots, c^{k+1}\epsilon_0).$$

This suggest the following norm in ϵ -space:

$$\|\epsilon\| := \max \left(|\epsilon_{k-1}|^{1/2}, \dots, |\epsilon_0|^{1/(k+1)} \right), \quad (3.1)$$

and the following equivalence relation in ϵ -space

$$\epsilon \simeq \epsilon' \iff \exists c \in \mathbb{R}^+, (\epsilon'_{k-1}, \dots, \epsilon'_0) = (c^2\epsilon_{k-1}, \dots, c^{k+1}\epsilon_0). \quad (3.2)$$

We call

$$\mathbb{D}_\rho = \{\epsilon; \|\epsilon\| \leq \rho\}.$$

Remark 3.1 We should normally describe the modulus on an open polydisk. We prefer to work with the closed polydisk \mathbb{D}_ρ , which allows to use its conic structure under the equivalence relation (3.2). Whenever we consider \mathbb{D}_ρ it is implicitly meant that the family is defined for ϵ in a larger open polydisk.

We consider the open component Σ_0 of \mathbb{D}_ρ which is the set of ϵ for which the discriminant $\Delta(\epsilon)$ of P_ϵ is nonzero, i.e. the vector field is nondegenerate. It is a union of curves $\{(c^2\epsilon_{k-1}, \dots, c^{k+1}\epsilon_0) \mid c \in (0, 1)\}$ for some ϵ with $\|\epsilon\| = \rho$, which are invariant under the equivalence relation \simeq .

Notation 3.2 We denote by T_C the translation by C :

$$T_C(Z) = Z + C.$$

4 The construction of Fatou coordinates

From now on we will limit ourselves to a prepared family

$$f_\epsilon(z) = z + P_\epsilon(z)(1 + Q_\epsilon(z) + P_\epsilon(z)R_\epsilon(z)). \quad (4.1)$$

Let us first discuss the general strategy. In general, there are obstructions to conjugate a germ of analytic family of the type (4.1) with its normal form. So the idea is to construct the conjugacies with the normal form on sectorial domains in z -space. These sectorial domains, there will be $2k$ of them, should be sufficiently large so that the conjugacies be almost unique. Then, because of the (almost) uniqueness, their comparison measures the obstruction to an analytic conjugacy with the normal form on a full neighborhood \mathbb{D}_r of the origin in z -space, and it is an analytic invariant for the system. But we also want the conjugacies to depend analytically of ϵ . So we want our sectors in z -space and the conjugacies defined on them to depend analytically of ϵ . There are obstructions to do that because the sectors are attached both to some fixed part of the boundary of \mathbb{D}_r and to some fixed points which may rotate

around each other when the parameters vary. Hence the construction can only be done analytically in ϵ for ϵ in some sectors in parameter space.

The analytic classification of analytic diffeomorphisms (or their unfoldings) is closely related to the embedding problem which consists in determining when an analytic diffeomorphism is embeddable in a flow, i.e. conjugate to the time one map of a vector field. Since the normal form of f_ϵ is the time one map of the 1-dimensional vector field v_ϵ , this is the case when the diffeomorphism or family of diffeomorphisms is conjugate to its normal form. A standard practice in the field is to change coordinate to a coordinate in which the normal form is just the translation by 1, which we call T_1 (this coordinate is nothing else than the time of the vector field). Hence a Fatou coordinate is just a diffeomorphism conjugating the system to T_1 . But, in practice we will change coordinate for f_ϵ also.

4.1 Construction of the sectors in z -space

We want to describe the dynamics of the family f_ϵ over a fixed neighborhood \mathbb{D}_r of the origin for all values of the parameters in a polydisk \mathbb{D}_ρ . The construction of the $2k$ sectors for the different values of the parameters in a sectorial domain in parameter space is very similar to the construction of the squid sectors in [17]: in particular, quantitative versions of the qualitative constructions described below appear there. We recall the construction for purpose of completeness. It was only done for values of ϵ inside Σ_0 , (i.e. the fixed points are distinct (regular)) with continuous limit at $\epsilon = 0$, but it can easily be extended to all values of the parameters. For $\epsilon \in \Sigma_0$ (resp. $\epsilon \notin \Sigma_0$), each sector in z -space is adherent to two regular fixed points (resp. two regular or singular fixed points non necessarily distinct), i.e. the closure of the sector contains the fixed points .

The formal normal form of f_ϵ is the time one map of the vector field v_ϵ . For small z and ϵ , this vector field is close to the simpler vector field

$$w_\epsilon = P_\epsilon(z) \frac{\partial}{\partial z}. \quad (4.2)$$

The phase portrait of w_ϵ has been extensively studied in the literature (see for instance [2]), and we will be brief with the details. It is organized by its pole at infinity (see Figure 1(a)), which has order $k-1$ (hence for $k=1$, infinity is a regular point.) The pole has $2k$ separatrices which generically land on a singular point. In z -space we choose r sufficiently small so that the dynamics of f_0 be as in Figure 1(b): this comes from the fact that germs of diffeomorphisms of the form $f_0(z) = z + z^{k+1} + o(z^{k+1})$ are topologically equivalent to the time one map of the flow of the vector field $z^{k+1} \frac{\partial}{\partial z}$ (see for instance Theorem 2 of [4]). We then choose ρ sufficiently small so that all fixed points of f_ϵ remain inside \mathbb{D}_r , and the dynamics of f_ϵ be as in Figure 1(c)) along the boundary. For that purpose, we will in particular require that

$$\rho < \frac{r^k}{2}, \quad (4.3)$$

a condition that will be assumed throughout the paper. We will allow to further restrict ρ .

To construct normalizing changes of coordinates, it is standard to change coordinate and to work in the time coordinate of the vector field w_ϵ . This change of coordinate is given by

$$Z = p_\epsilon^{-1}(z) = \int_\infty^z \frac{dz}{P_\epsilon(z)} \quad (4.4)$$

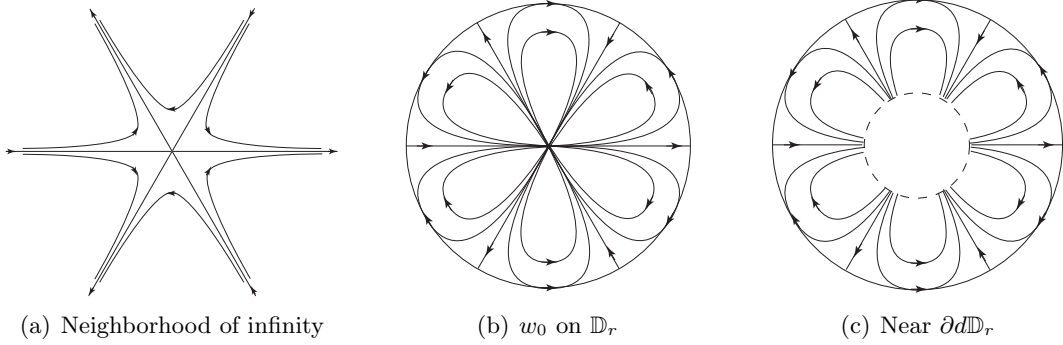


Figure 1: The phase portrait of w_ϵ at infinity and near the boundary of \mathbb{D}_r when $k = 3$.

In particular,

$$Z = p_\epsilon^{-1}(z) = \begin{cases} \sum_{j=1}^{k+1} \frac{1}{P'_\epsilon(z_j)} \ln(z - z_j), & z_j \text{ distinct,} \\ -\frac{1}{kz^k}, & \epsilon = 0. \end{cases}$$

In the Z -coordinate, the map f_ϵ is transformed into a map F_ϵ which is close to T_1 , the translation by 1:

Proposition 4.1 *The function F_ϵ is a small perturbation of the translation T_1 in the C^1 -topology. More precisely, there exists $K > 0$ such that, for $r > 0$ and $\rho > 0$ sufficiently small with ρ satisfying condition (4.3),*

$$|F_\epsilon(Z) - Z - 1| < Kr. \quad (4.5)$$

$$|F'_\epsilon(Z) - 1| < Kr^{k+1}. \quad (4.6)$$

PROOF. This kind of estimate is standard (see for instance [18]). For the sake of completeness we include the proof. On $\partial\mathbb{D}_r$ for small ϵ , p_ϵ^{-1} is a small perturbation of $-\frac{1}{kz^k}$. Hence (4.5) is satisfied for $|z| = r$ and ϵ sufficiently small. Since the function $(F_\epsilon(Z) - Z - 1) \circ p_\epsilon^{-1}$ is uniform in \mathbb{D}_r and bounded at the fixed points, by the maximum principle (4.5) is satisfied everywhere for the same values of ϵ .

Moreover, if we let $z = p_\epsilon(Z)$, then

$$F'_\epsilon(Z) = (p_\epsilon^{-1})'(f_\epsilon(z)) \cdot f'_\epsilon(z) \cdot p'_\epsilon(Z) = \frac{P_\epsilon(z)}{P_\epsilon(f_\epsilon(z))} f'_\epsilon(z).$$

Note that this function is analytic in (z, ϵ) since P_ϵ vanishes at the fixed points of f_ϵ . Let us now consider the case $\epsilon = 0$. Since $F'_0(Z) - 1 = \frac{P_0(z)f'_0(z) - P_0(f_0(z))}{P_0(f_0(z))}$, $P_0(f_0(z)) = z^{k+1} + (k+1)z^{2k+1} + O(z^{2k+2})$ and $P_0(z)f'_0(z) = z^{k+1} + (k+1)z^{2k+1} + O(z^{2k+2})$, then $F'_0(Z) - 1 = O(z^{k+1})$, yielding $|F'_0(Z) - 1| < \frac{K}{2}r^{k+1}$ for r sufficiently small. The estimate (4.6) will then be valid for $\|\epsilon\| < \rho$ with ρ sufficiently small. \square

The map p_ϵ is a k -sheeted covering of $\mathbb{C} \setminus \{z_1, \dots, z_{k+1}\}$. The inverse image of $\mathbb{CP}^1 \setminus \mathbb{D}_r$ is given by sequences of equidistant holes in the different sheets. Making a turn around one hole corresponds to making $1/k$ -th of a turn around $\partial\mathbb{D}_r$. Each outgoing separatrix of infinity

undergoes a homoclinic connection with an incoming separatrix of infinity when the vector field w_ϵ is transformed into a vector field $w_\epsilon(\theta) = e^{i\theta}w_\epsilon$ for some θ to be discussed later. The distance between the centers of two consecutive holes corresponds to $e^{-i\theta}T$, where $T \in \mathbb{R}^+$ is the travel time along the homoclinic loop of $w_\epsilon(\theta)$ (a pole is reached in finite time). This distance is also equal to $\pm \sum_{j \in I} \frac{2\pi i}{P'_\epsilon(x_j)}$, where $I \subset \{1, \dots, k+1\}$ is the set of indices of the singular points on one side of the homoclinic loop of $w_\epsilon(\theta)$. We will describe these distances in Section 4.2 after we have described the sectors in ϵ -space covering Σ_0 .

To simplify the presentation, it will suffice to construct the modulus for parameter values inside Σ_0 , (i.e. the fixed points are distinct). Note that $\Sigma_0 = \{\epsilon \mid \Delta(\epsilon) \neq 0\}$, where $\Delta(\epsilon)$ is the discriminant of P_ϵ . Since the construction will be bounded near the points of $\Delta = 0$, it will be possible to fill the holes later. The sectors will be bounded by four curves:

- (i) a portion S of $\partial\mathbb{D}_r$ of opening $\beta \in (\frac{\pi}{k}, \frac{3\pi}{2k})$,
- (ii) two “complex time trajectories” of w_ϵ starting at the end points of S and ending in two distinct singular points,
- (iii) a “complex time trajectory” of w_ϵ joining the two singular points inside \mathbb{D}_r .

The curves can be chosen so that any three sectors have a void intersection.

The idea is that the sectors should contain representatives of a *large* set of orbits of f_ϵ parameterized by \mathbb{C}^* . Since the linear maps are the only holomorphic diffeomorphisms of \mathbb{C}^* , this endows the set of orbits of f_ϵ over the sector of a conformal structure with a very rigid coordinate. To ensure that, it is best if the boundary of the sectors are transversal to curves invariant under f_ϵ . The simplest idea is to take for (ii) and (iii) trajectories of the orthogonal vector field iw_ϵ starting on $\partial\mathbb{D}_r$ and corresponding to vertical lines in the Z -coordinate. This was done by Oudkerk in [9]. However, for some values of ϵ for which iw_ϵ is not far from having a homoclinic loop, there may exist no such trajectories inside \mathbb{D}_r . In that case, we replace each trajectory of iw_ϵ by a union of two to three trajectories of $e^{i\theta}w_\epsilon$ for some $\theta \in [\frac{\pi}{4}, \frac{3\pi}{4}]$ (it could be different θ 's on the different pieces of trajectories.) Figure 2 (resp. Figure 3) represents such sectors in Z -space (resp. z -space).

When considering limits for $\epsilon \rightarrow 0$, it will suffice to consider such limits along the curves in parameter space that are invariant under the equivalence relation \simeq defined in (3.2). In the limit for $\epsilon = 0$, the sectors will be given by half planes in Z -space limited by vertical lines.

In practice, we want that when we consider ϵ in an invariant curve in parameter space, then for any compact set in the limit sector at $\epsilon = 0$, it is included in the corresponding sector for ϵ sufficiently small. For this purpose, each complex time trajectory of (ii) will be given by a finite piece of iw_ϵ starting on $\partial\mathbb{D}_r$ followed by a trajectory of $e^{i\theta}w_\epsilon$ ending in the singular point, for some $\theta \in [\frac{\pi}{4}, \frac{3\pi}{4}]$ (it could be different θ 's at the two singular points). The trajectory in (iii) is a union of trajectories of iw_ϵ and of the two $e^{i\theta}w_\epsilon$ used in the neighborhood of the singular points. It is the same θ that is used for all curves approaching a given singular point. Such sectors are best visualized in Z -space where they are bounded by vertical segments and segments, lines or half-lines making some angle $\theta \in [\frac{\pi}{4}, \frac{3\pi}{4}]$ with the horizontal direction (Figure 2).

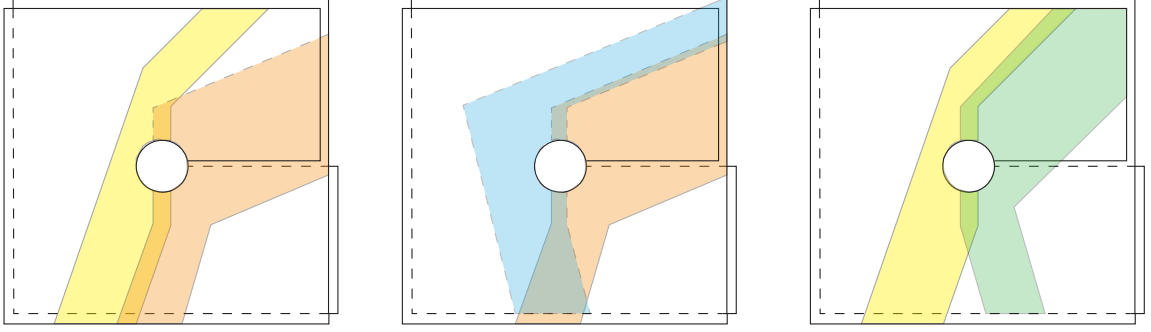


Figure 2: Different pairs of sets in Z -space corresponding to four sectors in z -space for the same value of ϵ .

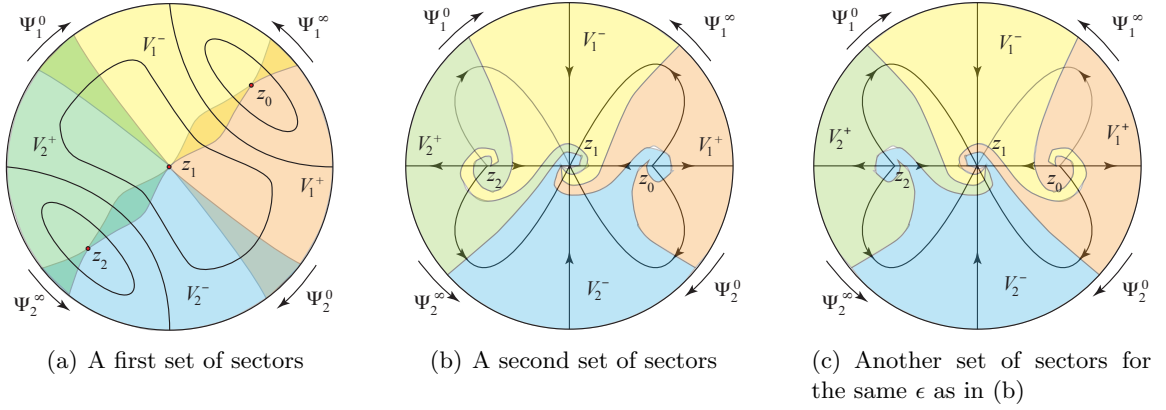


Figure 3: Different types of sets of sectors in z -space. In (b) and (c) their overlapping has not been drawn.

Remark 4.2 The construction can be carried also for values of $\epsilon \notin \Sigma_0$, for which $\Delta(\epsilon) = 0$. In that case, for some sectors the singular points of (ii) are identical and there is no trajectory of type (iii). In practice these sectors correspond to domains conformally equivalent to a half-plane in Z -space in a leaf where there is no sequence of holes (the distance between holes has become infinite).

Notation 4.3 The $2k$ sectors in z -space are called V_j^\pm (or $V_{j,\epsilon}^\pm$). The sectors V_j^+ (resp. V_j^-) are centered on regions of $\partial\mathbb{D}_r$ where f_ϵ exits (resp. enters) \mathbb{D}_r .

It remains to discuss how to choose the angles θ . For this, we need to make a digression in parameter space.

4.2 The sectors in parameter space

We first consider the case where $r = \infty$. Then, the construction described above can be done with $\theta = \pi/2$, except when iw_ϵ has a homoclinic trajectory: this occurs when there exists a nontrivial subset $I \subset \{1, \dots, k+1\}$ such that $\sum_{j \in I} \frac{i}{P'_\epsilon(x_j)} \in i\mathbb{R}$. The closure of this set in parameter space contains $\Delta = 0$. This can also be seen geometrically, since there is always an attracting and a repelling separatrix of infinity ending at a multiple point. The complement of the closure of this set in parameter space is a union of $C(k) = \frac{1}{k+1} \binom{2k}{k}$ connected components \widetilde{W}_ℓ . This comes from the following theorem which should be applied to iw_ϵ after a scaling in z .

Theorem 4.4 *In the parameter space of polynomial vector fields of the form $P_\epsilon(z) \frac{\partial}{\partial z}$ there exist exactly $C(k) = \frac{1}{k+1} \binom{2k}{k}$ different connected components of generic vector fields with simple singular points and no homoclinic loops. These components are simply connected.*

PROOF. This is easily derived from [2]. There, it is shown that there are exactly $C(k)$ combinatorial invariants describing how the separatrices of generic vector fields are attached to the singular points (see Figure 4). It suffices to show that each such combinatorial invariant corresponds to a simply connected open set in ϵ -space. Indeed, each vector field is characterized by its combinatorial invariant plus k analytic invariants $b_j \in \mathbb{C}$, such that $\text{Im } b_j > 0$, corresponding to the transit time from ∞ to ∞ along the dotted lines in Figure 4(b). These times are periods that can be calculated using the residue theorem. Each such vector field is realized on an abstract manifold M_ϵ constructed by sewing adequately $k+1$ horizontal strips, the width of which is given by the invariants b_j . Hence, the family M_ϵ , for ϵ in a connected component, is parameterized by $\mathcal{B} = \{\mathbf{b} = (b_1, \dots, b_k) \in \mathbb{H}^k\}$. This set is obviously simply connected. On a given connected component, we have a holomorphic bijection $\epsilon \mapsto \mathbf{b}$, yielding that the construction of M_ϵ depends holomorphically on $\mathbf{b} \in \mathcal{B}$. Hence, any connected component of generic vector fields in parameter space has the same topology as \mathbb{H}^k . \square

Sectors providing a covering of Σ_0 . The sectors $\{W_\ell \mid \ell = 1, \dots, C(k)\}$ in parameter space provide a covering of Σ_0 . Each sector W_ℓ has a vertex at $\epsilon = 0$, and is an enlargement of the corresponding \widetilde{W}_ℓ . Let us call Σ_1 the subset of Σ_0 for which iw_ϵ is generic, i.e. has no homoclinic trajectory. In practice, we will not be able to use the trajectories of iw_ϵ for all values of $\epsilon \in \widetilde{W}_\ell$. Indeed, it may happen that a separatrix of infinity makes large spirals disconnecting \mathbb{D}_r before landing at a singular point. If that is the case, then we replace part of corresponding trajectory of iw_ϵ by a trajectory of $e^{i\theta(\epsilon)}w_\epsilon$ for some $\theta(\epsilon)$ with

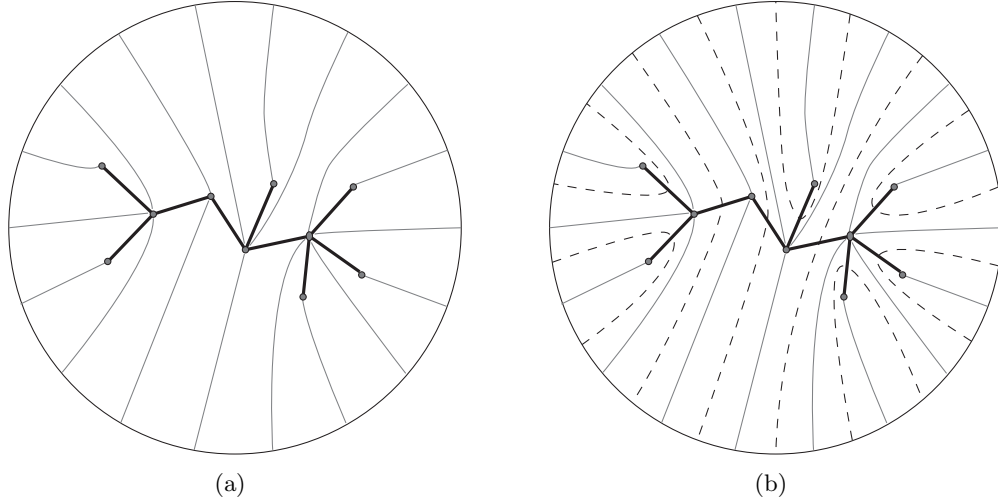


Figure 4: The graph (tree) formed by the separatrices of infinity ending at the singular points, together with trajectories joining the singular points. The distances between holes on the different sheets in Z -space are given by the complex travel times from ∞ to ∞ along the dotted lines in (b).

$\theta(\epsilon) \in (\frac{\pi}{4}, \frac{3\pi}{4})$ depending continuously on ϵ . This can be done not only for $\epsilon \in \widetilde{W}_\ell$, but we can also extend the \widetilde{W}_ℓ 's so that their union covers Σ_0 . Indeed, the boundaries of \widetilde{W}_ℓ consist of hypersurfaces where homoclinic connections occur. For each homoclinic connection, we choose an appropriate angle θ to destroy it. In particular, the sign of $\theta - \frac{\pi}{2}$ is such that the respective positions of the separatrices are the same as for \widetilde{W}_ℓ (this is ensured by the continuity of $\theta(\epsilon)$.) Each \widetilde{W}_ℓ contains a codimension k subset for which k homoclinic loops of w_ϵ (not iw_ϵ !) appear simultaneously. On this subset, which is the *organizing center* of W_ℓ , we use trajectories of iw_ϵ .

Remark 4.5 The distances of the holes are well defined for ϵ in some \widetilde{W}_ℓ : these distances are vectors. They are precisely $-i$ times the complex travel time along the trajectories in dotted lines of Figure 4(b) for the vector field iw_ϵ . There are k such times, one for each sheet in Z -space. For $\epsilon \in W_\ell$, we extend continuously these distances by considering $e^{-i\theta(\epsilon)}$ times the complex travel time along corresponding trajectories in dotted lines for the vector field $e^{i\theta(\epsilon)}w_\epsilon$, where $\theta(\epsilon)$ depends continuously on ϵ .

Note that for $\epsilon \in W_\ell \cap W_{\ell'}$ and $k > 2$ we may have distinct sets of distances depending whether we consider the W_ℓ or $W_{\ell'}$ point of view.

4.3 Construction of the Fatou coordinates

The construction is performed for $\epsilon \in W_\ell$.

Definition 4.6 We call a translation domain the saturation by F_ϵ of any preimage by p_ϵ of a sector in z as described in Section 4.1, provided that the Z -domain is sufficiently wide to contain a strip, $\hat{C}(\ell)$, bounded by a curve ℓ going from $\text{Im}(Z) = -\infty$ to $\text{Im}(Z) = +\infty$ with

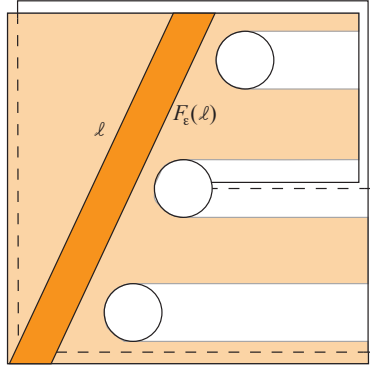


Figure 5: A translation domain.

slope greater than 1 at all points, and its image by F_ϵ (see Figure 5). The curve ℓ is called an admissible curve for the translation domain which we will denote by \mathcal{T}_ϵ .

The *Fatou coordinates* are changes of coordinates on the translation domains conjugating the diffeomorphism F_ϵ to the translation by 1, noted T_1 .

Theorem 4.7 *There exists $r > 0$ and $\rho > 0$ sufficiently small with $\rho < \frac{r^k}{2}$ so that for $\epsilon \in \mathbb{D}_\rho$, for \mathcal{T}_ϵ any translation domain, and $Z_0(\epsilon) \in \mathcal{T}_\epsilon$, then*

1. *there exists a holomorphic diffeomorphism $\Phi_\epsilon : \mathcal{T}_\epsilon \rightarrow \mathbb{C}$, such that*

$$\Phi_\epsilon \circ F_\epsilon = T_1 \circ \Phi_\epsilon, \quad (4.7)$$

for $Z \in \mathcal{T}_\epsilon \cap F_\epsilon^{-1}(\mathcal{T}_\epsilon)$. Moreover, $\lim_{\text{Im}(Z) \rightarrow \pm\infty} \text{Im}(\Phi_\epsilon(Z)) = \pm\infty$.

2. *If $\Phi_{1,\epsilon}$ and $\Phi_{2,\epsilon}$ are two solutions of (4.7), then there exists $A \in \mathbb{C}$ such that $\Phi_{2,\epsilon}(Z) = T_A \circ \Phi_{1,\epsilon}(Z)$. In particular, if $Z_0(\epsilon) \in \mathcal{T}_\epsilon$, there exists a unique holomorphic diffeomorphism Φ_ϵ satisfying (4.7) together with $\Phi_\epsilon(Z_0(\epsilon)) = 0$.*
3. *If $Z_0(\epsilon)$ depends analytically on ϵ with limit $Z_0(0)$ at $\epsilon = 0$, then the unique Fatou coordinate Φ_ϵ satisfying $\Phi_\epsilon(Z_0(\epsilon)) = 0$ depends analytically on $\epsilon \in W_\ell$. For $\epsilon \rightarrow 0$, it converges to Φ_0 , uniformly on compact sets in \mathcal{T}_0 .*

PROOF. The proof is exactly the same as in [7] since it relies only on Proposition 4.1. We put it here for the sake of completeness. The technique we use is identical to that of Shishikura [18], as adapted in [7]. It consists in constructing a quasi-conformal conjugacy of F_ϵ to T_1 , and then using Ahlfors-Bers theorem to transform it into a conformal conjugacy.

All along the proof we do not mention the ϵ -dependence. Let ℓ be an admissible curve for the translation domain \mathcal{T} , and $\hat{C}(\ell)$ the corresponding strip. Let $\beta(t)$, $t \in \mathbb{R}$, be a parametrization of ℓ such that $\text{Im}(\beta(t)) = t$. We define $h_1 : C_0 = \{(X, Y) \in \mathbb{R}^2 \mid 0 \leq X \leq 1\} \rightarrow \hat{C}(\ell)$ by:

$$h_1(X, Y) = (1 - X)(\beta(Y)) + XF_\epsilon(\beta(Y)). \quad (4.8)$$

Then

$$\begin{aligned}\frac{\partial h_1}{\partial X} &= F_\epsilon(\beta(Y)) - \beta(Y) \\ \frac{\partial h_1}{\partial Y} &= \beta'(Y) [1 + X(F'_\epsilon(\beta(Y)) - 1)].\end{aligned}\quad (4.9)$$

Using the estimates (4.5) and (4.6), these formulas imply that

$$\frac{\partial h_1}{\partial X} - 1 = u(X, Y), \quad \frac{\partial h_1}{\partial Y} - \beta'(Y) = v(X, Y), \quad \text{with } |u| < Kr, \quad |v| \leq 2Kr^{k+1}. \quad (4.10)$$

Let $\mu_{h_1} = \frac{\partial h_1}{\partial \bar{Z}} / \frac{\partial h_1}{\partial Z}$ be the dilatation coefficient field of h_1 . One has

$$\frac{\partial h_1}{\partial \bar{Z}} = \frac{1}{2}[1 + u + i(\beta'(Y) + v)] \quad \text{and} \quad \frac{\partial h_1}{\partial Z} = \frac{1}{2}[1 + u - i(\beta'(Y) + v)]. \quad (4.11)$$

Note that $\beta(Y) = \gamma(Y) + iY$, where $\gamma'(Y) \in [-1, 1]$. We have

$$\mu_{h_1} = \frac{u + iv + i\gamma'(Y)}{(2 - i\gamma'(Y)) + u - iv} < \frac{1 + Kr + Kr^{k+1}}{1 - Kr - Kr^{k+1}} < 1. \quad (4.12)$$

when r is sufficiently small, independently of ϵ such that $\|\epsilon\| \leq \rho < \frac{r^k}{2}$.

Hence, h_1 is a quasi-conformal mapping on the strip C_0 and satisfies $h_1^{-1}(F_\epsilon(Z)) = h_1^{-1}(Z) + 1$ for $Z \in \ell$. Moreover, $\mu = \mu_{h_1}$ is a Beltrami field on C_0 . (One can also write that $\mu = h_1^* \mu_0$, where μ_0 is the standard Beltrami field on \mathbb{C} (defined by the function 0).)

We extend μ to a Beltrami field on all of \mathbb{C} by means of the translation T_1 . Indeed, the extended μ is periodic of period 1, is in $L^\infty(\mathbb{C})$ and has a L^∞ -norm: $\|\mu\|_\infty = \|\mu_{h_1}\|_\infty < 1$.

The universal covering

$$w = E(W) = \exp(-2\pi i W) \quad (4.13)$$

from \mathbb{C} to \mathbb{C}^* induces a holomorphic diffeomorphism from \mathbb{C}/T_1 to \mathbb{C}^* . Since μ is invariant by T_1 , the map E induces a Beltrami field $\tilde{\mu}$ on \mathbb{C}^* with the same norm: $\mu = E^*(\tilde{\mu})$. We can extend $\tilde{\mu}$ to a Beltrami field on the Riemann sphere \mathbb{CP}^1 by letting, for instance, $\tilde{\mu}(0) = \tilde{\mu}(\infty) = 0$.

By Ahlfors-Bers measurable mapping theorem there exists a unique quasi-conformal mapping $\tilde{h}_2 : \mathbb{CP}^1 \rightarrow \mathbb{CP}^1$ such that $\tilde{h}_2^* \mu_0 = \tilde{\mu}$, and $\tilde{h}_2(0) = 0, \tilde{h}_2(\infty) = \infty, \tilde{h}_2(1) = 1$. Since $0, 1 \in E^{-1}(1)$, this map lifts to a quasi-conformal map $h_2 : \mathbb{C} \rightarrow \mathbb{C}$ sending 0 to 0, and 1 to 1. Indeed, one can lift \tilde{h}_2 into a map h_2 such that $h_2(0) = 0$. The circle in \mathbb{CP}^1 which turns one time around 0 or ∞ lifts into the line segment $[0, 1]$ in \mathbb{C} . This means that $h_2(1) = 1$. We have also that $\text{Im}(h_2(X + iY)) \rightarrow \pm\infty$ when $Y \rightarrow \pm\infty$.

Moreover, h_2 commutes with T_1 . Indeed, the homeomorphism $H_2 = h_2 \circ T_1 \circ h_2^{-1}$ induces the identity on \mathbb{CP}^1 and must then be a power of the deck transformation T_1 of the universal covering map E : i.e. $H_2 = T_1^n$ for some $n \in \mathbb{Z}$. Now $H_2(0) = h_2 \circ T_1(0) = h_2(1) = 1$. This forces $n = 1$, and then $H_2 = T_1$, i.e. $h_2 \circ T_1 = T_1 \circ h_2$.

We define $\phi : \hat{C}(\ell) \rightarrow \mathbb{C}$ by $\phi = h_2 \circ h_1^{-1}$, and extend it by T_1 to a mapping $\phi : \mathbb{C} \rightarrow \mathbb{C}$ which is quasi-conformal and preserves the standard conformal structure. Hence, it is conformal. For $Z \in \ell$, one has $T_1 \circ \phi(Z) = \phi \circ F_\epsilon(Z)$. Then ϕ extends to a map Φ of \mathcal{T} into \mathbb{C} by $\Phi(Z) = \phi \circ F_\epsilon^n(Z) - n$ where $n \in \mathbb{Z}$ is such that $F_\epsilon^n(Z) \in \hat{C}(\ell)$. This map Φ is a holomorphic diffeomorphism which verifies: $\Phi \circ F_\epsilon = T_1 \circ \Phi$.

If $\Phi_{i,\epsilon}$, $i = 1, 2$, satisfy (4.7), let $\chi_{1,\epsilon} = \Phi_{2,\epsilon} \circ (\Phi_{1,\epsilon})^{-1}$ and $\chi_{2,\epsilon} = \Phi_{1,\epsilon} \circ (\Phi_{2,\epsilon})^{-1}$. Both $\chi_{j,\epsilon}$ satisfy $\chi_{j,\epsilon}(W + 1) = \chi_{j,\epsilon}(W) + 1$ and this relation allows to extend them to global diffeomorphisms of \mathbb{C} commuting with T_1 . It follows that $\Phi_{2,\epsilon} \circ \Phi_{1,\epsilon}^{-1}$ is a translation.

By construction, it is standard that the Fatou coordinate Φ_ϵ depends analytically on ϵ , with uniform convergence to Φ_0 on compact sets on \mathcal{T}_0 when $\epsilon \rightarrow 0$. \square

4.4 Components of the analytic part of the modulus and normalization

Notation 4.8 We denote by $\mathcal{T}_{j,\epsilon}^\pm$ the translation domain containing $p_\epsilon^{-1}(V_{j,\epsilon}^\pm)$. We call $\Phi_{j,\epsilon}^\pm$ the Fatou coordinate on $\mathcal{T}_{j,\epsilon}^\pm$ containing $p_\epsilon^{-1}(V_{j,\epsilon}^\pm)$.

We have $2k$ strips given by $p_\epsilon^{-1}(V_{j,\epsilon}^\pm)$ inside the translation domains. These strips intersect. For $\epsilon \in \Sigma_0$, each strip intersects one or two of its neighbors along two semi-infinite strips. All together, this yields $2k$ semi-infinite strips. Let us call $\Phi_{j,\epsilon}^+$ (resp. $\Phi_{j,\epsilon}^-$) the Fatou coordinates corresponding to translation domains on the left (resp. right) of the principal hole as in Figure 2 (the Fatou coordinate $\Phi_{j,\epsilon}^\pm$ is defined on a translation domain $\mathcal{T}_{j,\epsilon}^\pm$ containing $p_\epsilon^{-1}(V_{j,\epsilon}^\pm)$).

Comparison of the Fatou coordinates. For this purpose we introduce the maps:

$$\begin{cases} \Psi_{j,\epsilon}^0 = \Phi_{j,\epsilon}^- \circ (\Phi_{j+1,\epsilon}^+)^{-1}, \\ \Psi_{j,\epsilon}^\infty = \Phi_{j,\epsilon}^- \circ (\Phi_{j,\epsilon}^+)^{-1}. \end{cases} \quad (4.14)$$

Note that there exists $Y_0 > 0$ such that $\Psi_{j,\epsilon}^0$ (resp. $\Psi_{j,\epsilon}^\infty$) is defined in a region $\text{Im}(Z) < Y_0$ (resp. $\text{Im}(Z) > Y_0$). These maps are almost the *analytic part of the modulus*, which will indeed almost be given by a collection of $2k$ families of analytic maps. But, the maps $\Psi_{j,\epsilon}^0$ and $\Psi_{j,\epsilon}^\infty$ are only defined up to composition with left and right translations. So the modulus will be given by the corresponding quotient. It is useful to work with representatives that depend nicely of the parameters, and in particular that have a limit when $\epsilon \rightarrow 0$ inside a sector W_ℓ . For that purpose, we introduce a normalization.

Lemma 4.9 *The transition maps $\Psi_{j,\epsilon}^{0,\infty}$, $j = 1, \dots, k$ can be expanded as Fourier series with constant terms $A_{j,\epsilon}^{0,\infty}$. It is possible to choose the base points of the Fatou coordinates in Theorem 4.7 so that, for all j , $A_{j,\epsilon}^0 = -A_{j,\epsilon}^\infty = \frac{\pi i a(\epsilon)}{k}$.*

PROOF. The proof of this fact is the same as in [11]. \square

This leaves one degree of freedom in the choice of one base point which can be chosen depending analytically on ϵ . We use it in the following way: we choose one base point $z_0 \in V_1^+ \subset \mathbb{D}_r$ which is a regular point: for instance, $z_0 = \frac{3r}{4}$. Let $Z_0(\epsilon) = p_\epsilon^{-1}(z_0) \in \mathcal{T}_{1,\epsilon}^+$. We normalize the first Fatou coordinate so that $\Phi_{1,\epsilon}^+(Z_0(\epsilon)) = 0$. Then the only freedom of choice is the base point $z_0 \in V_1^+$: it is now independent on ϵ .

Definition 4.10 *A set of Fatou coordinate is called a normalized set of Fatou coordinates if for all j , the constant terms $A_{j,\epsilon}^{0,\infty}$ of $\Psi_{j,\epsilon}^{0,\infty}$ satisfy $A_{j,\epsilon}^0 = -A_{j,\epsilon}^\infty = \frac{\pi i a(\epsilon)}{k}$. If, furthermore, there exists $z_0 \in \mathbb{D}_r$ which, for all $\epsilon \in \mathbb{D}_\rho$, is a regular point of f_ϵ inside $V_{1,\epsilon}^+$, and such that $\Phi_{1,\epsilon}^+(p_\epsilon^{-1}(z_0)) = 0$, then the set of Fatou coordinates is said to be strongly normalized.*

5 The theorem of analytic classification

Notation 5.1 For each sectorial domain W_ℓ in Σ_0 , we denote by $\Psi_{j,\epsilon,\ell}^{0,\infty}$ the elements of the corresponding analytic part of the modulus.

Theorem 5.2 *Two germs of generic prepared families f_ϵ and \tilde{f}_ϵ unfolding a germ of diffeomorphism with a parabolic point of codimension k are analytically conjugate if and only if they have the same normalized modulus, i.e. there exists $m \in \mathbb{Z}_k$ such that*

- for all $j = 0, \dots, k-1$, $\tilde{\epsilon}_j = \exp(-2\pi i m(j-1)/k)\epsilon_j$, and $a(\epsilon) = \tilde{a}(\tilde{\epsilon})$;
- for each each sectorial domain W_ℓ in Σ_0 , there exists $C_\ell(\epsilon)$ depending analytically on $\epsilon \in W_\ell$, with $C_\ell(0)$ independent of ℓ , such that for $j = 1, \dots, k$

$$\Psi_{j,\epsilon,\ell}^{0,\infty} = T_{-C_\ell(\epsilon)} \circ \tilde{\Psi}_{j+m,\epsilon,\ell}^{0,\infty} \circ T_{C_\ell(\epsilon)}.$$

In particular, a germ of generic prepared family f_ϵ unfolding a germ of diffeomorphism with a parabolic point of codimension k is analytically equivalent to its normal form if and only if, for all $j = 1, \dots, k$ and all $\epsilon \in \Sigma_0$, $\Psi_{j,\epsilon}^0 = T_{\frac{\pi i a(\epsilon)}{k}}$, $\Psi_{j,\epsilon}^\infty = T_{-\frac{\pi i a(\epsilon)}{k}}$.

PROOF. Suppose that the two prepared families f_ϵ and \tilde{f}_ϵ are analytically conjugate. The first condition (formal part) follows from Theorem 2.2. The canonical parameters are either the same, or such that $\tilde{\epsilon}_j = \exp(-2\pi i m(j-1)/k)\epsilon_j$ for some fixed m , and we can of course take ρ sufficiently small so the moduli of f_ϵ and \tilde{f}_ϵ be defined for $\epsilon \in \mathbb{D}_\rho$. Let h_ϵ be a conjugacy between f_ϵ and \tilde{f}_ϵ , such that $h_\epsilon \circ f_\epsilon = \tilde{f}_\epsilon \circ h_\epsilon$. Then $H_\epsilon = p_\epsilon \circ h_\epsilon \circ p_\epsilon^{-1}$ is a conjugacy between F_ϵ and \tilde{F}_ϵ , yielding that $\Phi_{j,\epsilon}^{0,\infty} \circ H_\epsilon^{-1}$ is a Fatou coordinate for \tilde{f}_ϵ such that $\Phi_{j,\epsilon}^{0,\infty} \circ H_\epsilon^{-1}(p_\epsilon^{-1}(z_0)) = 0$. The result follows by (almost) uniqueness properties of Fatou coordinates.

Conversely, consider two families with same normalized moduli. We can always conjugate the second family with a rotation and suppose that the two families have the same canonical parameter ϵ . Then, for $\epsilon \in W_\ell$, a conjugacy between the two families is given by

$$h_{\epsilon,\ell} = p_\epsilon \circ [(\tilde{\Phi}_{j,\epsilon,\ell}^\pm)^{-1} \circ T_{C_\ell(\epsilon)} \circ \Phi_{j,\epsilon,\ell}^\pm] \circ p_\epsilon^{-1}, \quad z \in V_{j,\epsilon,\ell}^\pm.$$

It is well defined because the two families have the same modulus, and it depends analytically on $\epsilon \in W_\ell$. We first show that in the generic case where the modulus is non trivial (i.e. at least one component $\Psi_{j,\epsilon}^{0,\infty}$ is not affine) the $h_{\epsilon,\ell}$ are independent of ℓ , and hence glue in a uniform h_ϵ on Σ_0 . Since, by construction, h_ϵ is bounded near $\{\Delta(\epsilon) = 0\}$, it can be extended to \mathbb{D}_ρ . We then show how to treat the non generic case.

So, let us try to show that $h_{\epsilon,\ell}$ and $h_{\epsilon,\ell'}$ coincide for $\epsilon \in W_\ell \cap W_{\ell'}$ and see what condition is necessary for this. We have that $g_\epsilon = (h_{\epsilon,\ell'})^{-1} \circ h_{\epsilon,\ell}$ is a symmetry of the first family: it commutes with f_ϵ . The families of diffeomorphisms commuting with f_ϵ are determined exactly as in the case $\epsilon = 0$, and the result is the same: they are given by powers $f_\epsilon^{\alpha(\epsilon)}$ of f_ϵ , where

- $\alpha(\epsilon) \in \mathbb{C}$, if all $\Psi_{j,\epsilon}^{0,\infty}$ are translations;

- $\alpha(\epsilon) \in \mathbb{Z}$, when the only translations T_b which commute with all $\Psi_{j,\epsilon}^{0,\infty}$ are the translations with $b \in \mathbb{Z}$;
- $\alpha(\epsilon) \in \frac{1}{n}\mathbb{Z}$, when the only translations T_b which commute with all $\Psi_{j,\epsilon}^{0,\infty}$ are the translations with $b \in \frac{1}{n}\mathbb{Z}$ for some $n \in \mathbb{N}^*$.

(Full details for the argument in the case $k = 1$ can be found in [14].) In the two last cases the symmetries are discrete, yielding that $\alpha(\epsilon)$ is constant. Since $C_\ell(0)$ is independent of ℓ and all $\Psi_{j,\epsilon,\ell}^{0,\infty}$ have the same limit $\Psi_{j,0}^{0,\infty}$ independent of ℓ , then $g_0 = id$, yielding that $\alpha(\epsilon) \equiv 0$.

In the non generic case of a trivial modulus, i.e. when all $\Psi_{j,\epsilon,\ell}^{0,\infty}$ are translations, then the modulus is invariant under conjugacy with a translation. So we can of course suppose that the modulus of f_ϵ and \tilde{f}_ϵ are strongly normalized. Also, because the components of the modulus are translations we can obviously take $C_\ell(\epsilon) \equiv 0$, which implies that $h_{\epsilon,\ell}(z_0) = h_{\epsilon,\ell'}(z_0) = \tilde{z}_0$ for all ℓ , i.e. $g_\epsilon(z_0) = z_0$. Then, necessarily $\alpha(\epsilon) \equiv 0$. \square

6 Geometric interpretation and parametric resurgence phenomenon

The codimension 1 case is very different from the codimension k case, where $k > 1$. Indeed, in the codimension 1 case, the analytic part of the modulus contains two functions, Ψ_ϵ^0 and Ψ_ϵ^∞ , and there are two fixed points. The generic case is the case where the two fixed points have multipliers with norm different from 1 and hence are linearizable, yielding that the system can be brought to the normal form in the neighborhood of each fixed point. In that case, the nontriviality of the modulus (i.e. at least one of Ψ_ϵ^0 or Ψ_ϵ^∞ is not a translation) expresses the fact that the two normalizing changes of coordinates do not match. Note that the domains of Ψ_ϵ^0 or Ψ_ϵ^∞ correspond to sectors in z -space attached to exactly one fixed point (see for instance [8] and [3]). So Ψ_ϵ^0 is attached to one fixed point, and Ψ_ϵ^∞ to the other. When the norm of the multiplier at one fixed point is equal to 1, then the linearizability or non linearizability of the fixed point can be determined from the Ψ_ϵ^0 or Ψ_ϵ^∞ which is attached to it. The parametric resurgence phenomenon is the phenomenon where the nonlinearizability of a fixed point for sequences of resonant values of ϵ (i.e. values of ϵ such that the multiplier at one fixed point is a root of unity) converging to the origin can be decided directly from Ψ_0^0 or Ψ_0^∞ . We will explain it in more generality below.

When $k > 1$, then $k + 1 \neq 2k$, and we lose the fact that exactly one component $\Psi_{j,\epsilon}^0$ or $\Psi_{j,\epsilon}^\infty$ of the modulus is attached to each fixed point. So the $2k$ components of the modulus should rather be considered as attached to the $2k$ sectors on the boundary, when we normalize the diffeomorphism from the boundary of $\partial\mathbb{D}_r$. What remains is that the $\Psi_{j,\epsilon}^0$ and $\Psi_{j,\epsilon}^\infty$ are attached to the fixed points in a surjective way. Moreover, all components attached to a fixed point are of the same type, either all $\Psi_{j,\epsilon}^0$, or all $\Psi_{j,\epsilon}^\infty$.

In Section 4.4 we have discussed the fact that the strips given by $p_\epsilon^{-1}(V_{j,\epsilon}^\pm)$ intersect along half strip(s). But in Figure 3 we also see some intersection parts that go from one fixed point to another. These are given by the projections of infinite strips in Z -space. These infinite strips are only obtained when one considers the periodicity in Z -space. The k different periods are given by the k different distances between holes. Then, an infinite strip is obtained by

intersection of one strip with the translation by a period of a second strip. On an infinite strip, the comparison of two Fatou coordinates is just a translation $T_{L(\epsilon)}(W) = W + L(\epsilon)$, called a *Lavaurs translation*, with $L(\epsilon) \rightarrow \infty$ when $\epsilon \rightarrow 0$. This comes from the fact that the comparison is a diffeomorphism commuting with T_1 on an infinite strip from $\text{Im}(W) = -\infty$ to $\text{Im}(W) = \infty$.

Now, to describe the dynamics near a fixed point, we use the renormalized return map: this map is defined on the orbit space, or its covering space, namely the Z -variable. It is the composition of the transition maps encountered when making one turn in the positive direction around the fixed point. These transition maps are the $\Psi_{j,\epsilon}^0$ or the $\Psi_{j,\epsilon}^\infty$ attached to the fixed point, alternating with as many Lavaurs translations.

It is easier to visualize the renormalized return maps when changing to the coordinate

$$w = \exp(-2\pi i Z).$$

The maps $\Psi_{j,\epsilon}^{0,\infty}$ expressed in this coordinate will be denoted $\psi_{j,\epsilon}^{0,\infty}$. The Lavaurs translations become Lavaurs linear maps $L_{c(\epsilon)}(w) = c(\epsilon)w$. Then the renormalized return map has a fixed point at 0 (resp. ∞) when built with $\psi_{j,\epsilon}^0$ (resp. $\psi_{j,\epsilon}^\infty$) and Lavaurs linear maps.

Typically a renormalized return in w -coordinate is of one of the forms

$$\begin{cases} L_{c_s(\epsilon)} \circ \psi_{j_s,\epsilon}^0 \circ \dots \circ L_{c_1(\epsilon)} \circ \psi_{j_1,\epsilon}^0, \\ L_{c_s(\epsilon)} \circ \psi_{j_s,\epsilon}^\infty \circ \dots \circ L_{c_1(\epsilon)} \circ \psi_{j_1,\epsilon}^\infty, \\ \psi_{j_s,\epsilon}^0 \circ L_{c_s(\epsilon)} \circ \dots \circ \psi_{j_1,\epsilon}^0 \circ L_{c_1(\epsilon)}, \\ \psi_{j_s,\epsilon}^\infty \circ L_{c_s(\epsilon)} \circ \dots \circ \psi_{j_1,\epsilon}^\infty \circ L_{c_1(\epsilon)}. \end{cases}$$

If $\lambda = \exp(\mu)$ is the multiplier of f_ϵ at one fixed point, then the multiplier at the fixed point of the renormalized return map in w coordinate is simply $\lambda' = \exp\left(\frac{4\pi^2}{\mu}\right)$.

The *parametric resurgence phenomenon* occurs when λ' is resonant: $\lambda' = \exp\left(2\pi i \frac{p}{q}\right)$. Note that this can occur along sequences $\{\epsilon_n\}$ of parameter values converging to 0, corresponding for instance to $\mu_n = -2\pi i \frac{q}{p+nq}$, which are located inside an open sector W_ℓ , and for which the constant $c(\epsilon_n)$ in the Lavaurs linear maps have a limit $c(0)$ when $n \rightarrow \infty$. Then, on such a sequence $\{\epsilon_n\}$ of parameter values, the coefficients of the normal form converge to the coefficients of the normal form of

$$\begin{cases} L_{c_s(0)} \circ \psi_{j_s,0}^0 \circ \dots \circ L_{c_1(0)} \circ \psi_{j_1,0}^0, \\ L_{c_s(0)} \circ \psi_{j_s,0}^\infty \circ \dots \circ L_{c_1(0)} \circ \psi_{j_1,0}^\infty, \\ \psi_{j_s,0}^0 \circ L_{c_s(0)} \circ \dots \circ \psi_{j_1,0}^0 \circ L_{c_1(0)}, \\ \psi_{j_s,0}^\infty \circ L_{c_s(0)} \circ \dots \circ \psi_{j_1,0}^\infty \circ L_{c_1(0)}. \end{cases}$$

This implies that if one of these maps is nonlinearizable of order m (because of the non vanishing of the m -th nonlinear term in the normal form) then, for n sufficiently large, the corresponding renormalized map is nonlinearizable of order $m'(\epsilon_n) \leq m$. As a general rule, the parametric resurgence phenomenon concerns cases where the only way that there be some incompatibility between the different sectorial normal forms on the different $V_{j,\epsilon}^\pm$ is that the fixed points themselves be non linearizable.

Applications in conformal geometry in the case $k = 1$ appear in [15] and [16].

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