# Germs of analytic families of diffeomorphisms unfolding a parabolic point (II) 

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## Structure of the mini-course

- Statement of the problem (first lecture)
- The preparation of the family (first lecture)
- Construction of a modulus of analytic classification in the codimension 1 case (second lecture)
- The realization problem in the codimension 1 case (third lecture)


## The parabolic germ



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Comparing is constructing a change of coordinates to normal form.

A global analytic comparison on $\mathbb{D}_{r}$ does not exist

So we cover $\mathbb{D}_{r}$ with two sectors $U_{ \pm}$. Over each sector the comparison is almost unique (up to a symmetry of the model, which is a time $t$ map of $v_{\epsilon}$.)

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The modulus is given by the comparison of the two normalizations over $U_{\cap}=U_{+} \cap U_{-}$. It is a symmetry of the model. (Because $U_{\cap}$ is small, there are many more symmetries .)

## The choice of the sectors $U_{ \pm}$



## The underlying idea



The dynamics is transversal to the inner part of the boundary of the sectors (except at the fixed point).

- It goes from $U_{+}$to $U_{-}$ on the part not joining the fixed points.
- It goes from $U_{-}$to $U_{+}$ on the part joining the fixed points.


## The choice of the sector in $\epsilon$

We work with a ramified covering of a neighborhood $\mathbb{D}_{\rho}$ of the origin in $\epsilon$-space.

$$
V_{\delta}=\{\hat{\epsilon}:|\hat{\epsilon}|<\rho, \arg \hat{\epsilon} \in(-\pi+\delta, 3 \pi-\delta)\}
$$

for $\delta \in(0, \pi)$.

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for $\delta \in(0, \pi)$.

$\delta$ can be taken arbitrarily small. The smaller $\delta$, the smaller $\rho$. In particular, the opening is always smaller than $4 \pi$.

The lift to the time coordinate

In practice we prefer to work with the time coordinate $Z$ of the vector field

$$
w_{\epsilon}=\left(z^{2}-\epsilon\right) \frac{\partial}{\partial z}
$$

namely we make the multivalued change of coordinate

$$
Z=p_{\epsilon}^{-1}(z)= \begin{cases}\frac{1}{2 \sqrt{\epsilon}} \ln \frac{z-\sqrt{\epsilon}}{z+\sqrt{\epsilon}} & \epsilon \neq 0 \\ -\frac{1}{z} & \epsilon=0\end{cases}
$$

with period $\alpha=\frac{\pi i}{\sqrt{\widehat{\epsilon}}}$ when $\epsilon \neq 0$.

## The underlying idea

If we allow complex time, then all points $z \in \mathbb{D}_{r}$ are in the trajectory of a unique point $z_{0}$. Hence, $z=\phi_{w_{0}}^{Z}\left(z_{0}\right)$ for some $Z$.

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The boundaries of the spiral sectors come from lines in Z-space.

## The sectors in Z-space



## The Fatou coordinates

Let $\Omega_{\widehat{\hat{E}}}^{ \pm}$be the sectors in $Z$-space. We construct Fatou coordinates $\Phi_{\hat{\epsilon}}^{ \pm}: \Omega_{\hat{\epsilon}}^{ \pm} \rightarrow \mathbb{C}$ which conjugate $F_{\epsilon}$ (the lifting of $f_{\epsilon}$ in $Z$-coordinate) to $T_{1}$ the translation by 1 :

$$
\Phi_{\hat{\epsilon}}^{ \pm} \circ F_{\epsilon}=T_{1} \circ \Phi_{\widehat{\epsilon}}
$$

## The construction of Fatou coordinates

We take a slanted line $\ell$ in $\Omega_{\widehat{\epsilon}}^{ \pm}$so that $\ell \cap F_{\epsilon}(\ell)=$ $\emptyset$, and so that the strip $S$ between $\ell$ and $F(\ell)$ is included in $\Omega_{\hat{\epsilon}}^{ \pm}$.


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We construct $\phi: S \rightarrow \mathbb{C}$ conjugating $F_{\epsilon}$ with $T_{1}$ by linear interpolation. We extend $\phi$. The map $\phi$ is quasi-conformal. We correct $\phi$ to a conformal map by Ahlfors-Bers theorem.

## The strips are fundamental domains in z-space



## Dependence of the Fatou coordinates on $\widehat{\epsilon}$

We can construct the Fatou coordinates $\Phi_{\hat{\hat{\epsilon}}}^{ \pm}$so that they depend analytically on $\hat{\epsilon}$ with continuous limit at $\hat{\epsilon}=0$.

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are unique modulo left composition with a translation.
In particular, they are uniquely defined by a base point such that

$$
\Phi_{\hat{\epsilon}}^{ \pm}(Z(\hat{\epsilon}))=0
$$

It suffices to take $Z(\hat{\varepsilon})$ depending analytically on $\hat{\varepsilon}$ with continuous limit at $\hat{\epsilon}=0$.

## The modulus of analytic classification

The sectors $\Omega_{\widehat{\widehat{\epsilon}}}^{ \pm}$intersect along two strips $\Omega_{\hat{\varepsilon}}^{0}$ and $\Omega_{\hat{\varepsilon}}^{\infty}$ where we can compare the Fatou coordinates


## The modulus of analytic classification

We define

$$
\left\{\begin{array}{lll}
\Psi_{\hat{\epsilon}}^{\infty}=\Phi_{\widehat{\hat{E}}}^{-} \circ\left(\Phi_{\hat{e}}^{+}\right)^{-1} & \text { on } \Omega_{\hat{e}}^{\infty} \\
\Psi_{\hat{E}}^{0}=\Phi_{\hat{\epsilon}}^{-} \circ\left(\Phi_{\hat{\epsilon}}^{+}\right)^{-1} & \text { on } \Omega_{\hat{E}}^{0}
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\end{array}\right.
$$

The modulus is defined as the equivalence class of

$$
\left(\Psi_{\hat{\epsilon}}^{0}, \Psi_{\hat{\epsilon}}^{\infty}\right)_{\hat{\epsilon} \in V_{\delta}}
$$



## The equivalence relation

$$
\left(\Psi_{\widehat{\epsilon}}^{0}, \Psi_{\widehat{\epsilon}}^{\infty}\right)_{\hat{\epsilon} \in V_{\delta}} \sim\left(\check{\Psi}_{\hat{\epsilon}}^{0}, \check{\Psi}_{\widehat{\epsilon}}^{\infty}\right)_{\hat{\epsilon} \in V_{\delta}}
$$

if and only if there exists $C_{\hat{\epsilon}}$ and $C_{\hat{\epsilon}}^{\prime}$ depending analytically on $\hat{\epsilon}$ with continuous limit at $\widehat{\epsilon}=0$ such that

$$
\left\{\begin{array}{l}
\Psi_{\widehat{\epsilon}}^{0}=T_{C_{\widehat{\epsilon}}} \circ \check{\Psi}_{\widehat{\epsilon}}^{0} \circ T_{C_{\hat{e}}^{\prime}} \\
\Psi_{\widehat{\epsilon}}^{\infty}=T_{C_{\widehat{\varepsilon}}} \circ \check{\Psi}_{\widehat{\epsilon}}^{\infty} \circ T_{C_{\widehat{\epsilon}}^{\prime}}
\end{array}\right.
$$

Theorem. [MRR] Two germs of generic families unfolding a codimension 1 parabolic point are analytically conjugate if and only if they have the same formal invariant $a(\epsilon)$ and the same modulus

$$
\left[\left(\Psi_{\hat{\epsilon}}^{0}, \Psi_{\hat{\epsilon}}^{\infty}\right)_{\hat{\epsilon} \in V_{\delta}}\right] / \sim
$$

## The proof

Let us suppose that two germs of prepared diffeomorphisms $f_{\epsilon}$ and $\tilde{f}_{\epsilon}$ have the same modulus. We can of course adjust the Fatou coordinates $\Phi_{\hat{\epsilon}}^{ \pm}$and $\widetilde{\Phi}_{\hat{\epsilon}}^{ \pm}$so that the representatives of the modulus be the same:

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\Psi_{\widehat{\epsilon}}^{0, \infty}=\widetilde{\Psi}_{\widehat{\epsilon}}^{0, \infty}
$$

Then a conjugacy $g_{\hat{\epsilon}}$ between the two systems is given by

$$
g_{\hat{\epsilon}}=\left\{\begin{array}{lll}
p_{\epsilon} \circ\left(\widetilde{\Phi}_{\hat{\epsilon}}^{+}\right)^{-1} \circ \Phi_{\hat{\epsilon}}^{+} \circ p_{\epsilon}^{-1} & \text { on } & U_{+} \\
p_{\epsilon} \circ\left(\widetilde{\Phi}_{\hat{\epsilon}}^{-}\right)^{-1} \circ \Phi_{\hat{\epsilon}}^{-} \circ p_{\epsilon}^{-1} & \text { on } & U_{-}
\end{array}\right.
$$

## Correction to a uniform conjugacy

We consider values

$$
\left\{\begin{array}{l}
\bar{\epsilon}=\hat{\epsilon} \\
\check{\epsilon}=\widehat{\epsilon} e^{2 \pi i}
\end{array}\right.
$$

Then

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h_{\epsilon}=\check{g}_{\epsilon}^{-1} \circ \bar{g}_{\epsilon}
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If $f_{\epsilon}$ has few symmetries then we can deduce that $h_{\epsilon} \equiv i d$. Otherwise, we need to correct the family of conjugacies $g_{\hat{e}}$ to a uniform family.

## The symmetries of $f_{\epsilon}$

We read them in the $W=\Phi(Z)$ variable. In this variable the dynamics is given by $T_{1}$. Hence the symmetries on the image of a Fatou coordinate are given by translations.

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To be global these symmetries need to commute with the modulus components $\Psi_{\hat{\epsilon}}^{0}$ and $\Psi_{\hat{\kappa}}^{\infty}$.

$$
\left\{\begin{array}{l}
\Psi_{\hat{\epsilon}}^{0}(W)=W+\sum_{n<0} b_{n}(\hat{\epsilon}) \exp (2 \pi i n W) \\
\Psi_{\widehat{\epsilon}}^{\infty}(W)=W-2 \pi i a(\epsilon)+\sum_{n>0} c_{n}(\widehat{\epsilon}) \exp (2 \pi i n W)
\end{array}\right.
$$

## Two cases

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\end{array}\right.
$$

1. If one of $b_{n}$ or $c_{n}$ is not identically zero, then the symmetries are discrete (either the identity, or of the form $f_{\epsilon}^{\circ \frac{p}{q}}$ for some fixed $q$ independent of $\epsilon$ ). Since $h_{0}=i d$, then $h_{\epsilon}=i d$.
2. If $b_{n} \equiv 0$ and $c_{n} \equiv 0$ for all $n$, then all symmetries are of the form $f_{\epsilon}^{\circ t(\epsilon)}$ for $t(\epsilon) \in \mathbb{C}$. Hence

$$
h_{\epsilon}=f_{\epsilon}^{\circ t(\epsilon)}
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$$

We correct $g_{\hat{e}}$ to

$$
g_{\widehat{\epsilon}} \circ f_{\epsilon}^{\circ \tau(\epsilon)}
$$

such that

$$
\tau(\check{\epsilon})-\tau(\bar{\epsilon})=t(\epsilon)
$$

The parametric resurgence phenomenon

We prefer to present the modulus under the form

$$
\left(\psi_{\hat{\epsilon}}^{0}, \psi_{\hat{€}}^{\infty}\right)
$$

where

$$
\psi_{\overparen{\epsilon}}^{0, \infty}=E \circ \Psi_{\overparen{\epsilon}}^{0, \infty} \circ E^{-1}
$$

and

$$
E=\exp (-2 \pi i W) .
$$

Then

$$
\left\{\begin{array}{l}
\psi_{\hat{e}}^{0}:(\mathbb{C}, 0) \rightarrow(\mathbb{C}, 0) \\
\psi_{\hat{\epsilon}}^{\infty}:(\mathbb{C}, \infty) \rightarrow(\mathbb{C}, \infty)
\end{array}\right.
$$

We consider values of $\epsilon$ such that $f_{\epsilon}^{\prime}( \pm \sqrt{\epsilon}) \in \mathbb{S}^{1}$.

The renormalized return maps become

$$
\left\{\begin{array}{l}
h_{\epsilon}^{0}=\psi_{\epsilon}^{0} \circ L_{\epsilon} \\
h_{\epsilon}^{\infty}=\psi_{\epsilon}^{\infty} \circ L_{\epsilon}
\end{array}\right.
$$

We consider sequences $\left\{\epsilon_{n}\right\}$ of values of $\epsilon$ such that $\epsilon_{n} \rightarrow 0$ and

$$
L_{\epsilon_{n}}(w)=\exp \left(\frac{2 \pi i p}{q}\right) w
$$

## Parametric resurgence phenomenon

Then for sufficiently large $n$

- $h_{\epsilon_{n}}^{0}$ is non linearizable as soon as $\psi_{0}^{0} \circ L_{\epsilon_{n}}$ is not linearizable.
- As a consequence $f_{\epsilon_{n}}$ is non linearizable at $-\sqrt{\epsilon}$ as soon as $\psi_{0}^{0} \circ L_{\epsilon_{n}}$ is not linearizable.


## At the other singular point

For sequences $\left\{\epsilon_{n}\right\}$ such that $\epsilon_{n} \rightarrow 0$ and

$$
L_{\epsilon_{n}}(w) \exp \left(-4 \pi^{2} a(\epsilon)\right)=\exp \left(\frac{2 \pi i p}{q}\right) w
$$

Then for sufficiently large $n$

- $h_{\epsilon_{n}}^{\infty}$ is non linearizable as soon as $\psi_{0}^{\infty} \circ L_{\epsilon_{n}}$ is not linearizable.
- As a consequence $f_{\epsilon_{n}}$ is non linearizable at $\sqrt{\epsilon}$ as soon as $\psi_{\infty}^{\infty} \circ L_{\varepsilon_{n}}$ is not linearizable.


## Interpretation of $\left(\psi_{\epsilon}^{0}, \psi_{\epsilon}^{\infty}\right)$

For values of the multiplier on the unit circle, $\psi_{\epsilon}^{0}$ controls the dynamics at $-\sqrt{\epsilon}$ and $\psi_{\epsilon}^{\infty}$ at $+\sqrt{\epsilon}$.

Interpretation of $\left(\psi_{\epsilon}^{0}, \psi_{\epsilon}^{\infty}\right)$

For values of the multiplier on the unit circle, $\psi_{\epsilon}^{0}$ controls the dynamics at $-\sqrt{\epsilon}$ and $\psi_{\epsilon}^{\infty}$ at $+\sqrt{\epsilon}$.

Understanding the dependence of $\left(\psi_{\epsilon}^{0}, \psi_{\epsilon}^{\infty}\right)$ would allow to understand the dynamics of points whose multiplier corresponds to an irrational rotation.

## The codimension $k$ case

The strategy

- Define a modulus for generic values of the parameters for which all fixed points are distinct.


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- The generic values of the parameters in a neighborhood of the origin belong to a finite union of open sets $V_{j}$, all adherent to the origin in parameter space.


## The codimension $k$ case

The strategy

- Define a modulus for generic values of the parameters for which all fixed points are distinct.
- The generic values of the parameters in a neighborhood of the origin belong to a finite union of open sets $V_{j}$, all adherent to the origin in parameter space.
- Give a description of the modulus for values of the parameters in each $V_{j}$ which depends analytically on the parameters with continuous limit when $\epsilon \rightarrow 0$.


## This yields a complete modulus of analytic classification

We consider two germs of prepared families of diffeomorphisms with same modulus.

- They are analytically conjugate over each $V_{j}$.


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- The conjugacies are bounded when approaching codimension 1 parameter values (one double fixed point), so they can be extended to this case.


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- The conjugacies are bounded when approaching codimension 1 parameter values (one double fixed point), so they can be extended to this case.
- We fill the holes by Hartogs' theorem.

